

The creep–rupture properties and the initiation and growth of the grain-boundary cracks in the cobalt-base HS-21 alloy

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The effect of the grain-boundary microstructures on the creep–rupture properties and the initiation and growth of the grain-boundary cracks was investigated using four kinds of specimen of various grain-boundary microstructures in the cobalt-base HS-21 alloy at 1089 K in air. Both the rupture strength and the creep ductility increased with increasing mean value of the fractal dimension of the grain boundaries, D_{gb} . The strain to crack initiation was largest in the specimen of the highest D_{gb} value (1.241), while the strain was much the same in the specimens of the D_{gb} value less than 1.162. This was explained by the local variation in the grain-boundary microstructures in these specimens. The mean value of the fractal dimension of the grain-boundary fracture, D_f , was close to the value of D_{gb} , although the value of D_f was a little higher than that of D_{gb} in the specimens of the lower D_{gb} values. The fracture appearance changed from a brittle grain-boundary fracture to a ductile one with increasing values of D_{gb} and D_f . The crack-growth rate in the surface-notched specimens decreased with increasing value of D_{gb} . The threshold stress intensity factor for crack growth was higher in the specimens with the higher D_{gb} values.

1. Introduction

It is known that serrated grain boundaries are effective in improving the creep–rupture properties in many heat-resistant steels and alloys [1–8]. As Ashby and co-workers [9–11] revealed, serrated grain boundaries composed of grain-boundary ledges, bumps and curvature, retard the grain-boundary sliding, which leads to the initiation and growth of the grain-boundary microcracks at high temperatures. Yamamoto *et al.* [5] and Tanaka and co-workers [6–8] have also shown that the grain-boundary fracture surface of the specimen with serrated grain boundaries is a ductile fracture surface composed of grain-boundary steps, ledges and small dimples, while the fracture surface of the specimen with straight grain boundaries is a typical brittle grain-boundary fracture surface. However, it has been difficult, until recently, to describe quantitatively the grain-boundary microstructures and fracture surfaces because of their complexity.

Mandelbrot *et al.* [12] first revealed the fractal nature (self-similarity) of fracture surfaces in the impact-loaded and fractured specimens of steels. As Milman *et al.* [13] have recently reviewed, the concept of the fractal geometry developed by Mandelbrot [14] as well as the measurement techniques of the fractal dimension have been applied to the morphological analysis of fracture surfaces in metallic materials, glasses and ceramics. The concept of the fractal geometry has been applied [15, 16] to creep fracture and it was found that there is a correlation between the

fractal dimension of the grain-boundary fracture (the fractal dimension of the grain-boundary fracture surface profile in the two-dimensional section, D'_f , $1 < D'_f < 2$) and the creep–rupture properties in the heat-resistant alloys. The creep–rupture properties can also be correlated to the fractal dimension of the grain boundaries (the fractal dimension of the grain-boundary surface profile in the two-dimensional section, D'_{gb} , $1 < D'_{gb} < 2$) [17].

Gokhale and co-workers [18, 19] have discussed the experimental techniques for the estimation of the fractal dimension of fracture surfaces. They have also pointed out as a result of a statistical analysis that thorough profile sampling is required to achieve precision in determining the fractal dimension for inhomogeneous surfaces [19]. However, the local variation in the microstructures or in the fracture surface morphologies can be correlated to critical phenomena like the fracture initiation in the specimens, while the overall fracture process may be governed by the averaged microstructures.

The initiation and growth of the grain-boundary cracks, which is related to the creep–rupture properties, may depend on the ruggedness of grain boundaries or fracture surface structures in the materials. In this study, the effect of the grain-boundary microstructures on the creep–rupture properties and the initiation and growth of the grain-boundary cracks was investigated using four kinds of specimen with various grain-boundary morphologies in the cobalt-base

HS-21 alloy at 1089 K in air. The mean values of the fractal dimensions of the grain boundaries and of the grain-boundary fracture and their standard deviations were estimated in the scale range of the fractal analysis less than one grain-boundary length by the box-counting method [14, 20, 21]. The relationship between the creep–rupture properties, the initiation and growth of the grain-boundary cracks and these fractal dimensions in the HS-21 alloy was then discussed.

2. Experimental procedure

2.1. Materials and creep and creep–rupture experiments

Table I shows the chemical composition of the commercial wrought cobalt-base HS-21 alloy. The alloy bars of 16 mm diameter were heat-treated to obtain four kinds of specimen with different grain-boundary configuration while keeping the strength of the matrix almost constant. Table II shows the heat treatments, the mean value, D_{gb} , and the standard deviation, s_{gb} , of the fractal dimension of the grain boundaries, and the matrix hardness in the specimens of the HS-21 alloy. The grain diameter of these specimens, d_g , was 130 μm . The fractal dimension of the grain boundaries is defined as that of the grain-boundary surface profile in the two-dimensional section ($1 < D_{gb} < 2$). As will be shown later, the values of D_{gb} and s_{gb} were estimated on five grain boundaries in each specimen by the box-counting method in this study [14, 20, 21]. The range of fractal dimension is also shown in Table II.

Fig. 1 shows the optical micrographs of the heat-treated specimens of the HS-21 alloy. Large M_{23}C_6 type carbide particles are visible in the specimens of the higher D_{gb} values (Fig. 1c and d), while fine M_{23}C_6 precipitates can be observed in the grains (Fig. 1a–d [8]). Grain boundaries are considerably serrated in specimens of the higher D_{gb} values (Fig. 1c and d),

whereas they are almost straight in specimens of lower D_{gb} values. The heat-treated specimens were machined into creep specimens as shown in Fig. 2. Smooth, round, bar specimens of 5 mm diameter, d , and 30 mm gauge length were used for creep–rupture experiments (Fig. 2a), and those with a shallow transverse notch were employed for crack-growth experiments (Fig. 2b). Creep experiments were carried out using single lever-type creep–rupture equipment under the initial creep stresses of 88.2–176 MPa at 1089 K. Conditions of the creep–rupture experiment were chosen to cause grain-boundary fracture in all the specimens. The time to crack initiation was also examined on the smooth, round, bar specimens.

The surface crack length, l , of the surface-notched specimen was measured by examining the specimen surface using an optical microscope after interrupting creep tests and cooling the specimens to room temperature. The specimens were heated to the test temperature and were loaded again after the crack-length measurement. The creep stress of the notched specimens is shown by the gross stress, σ_g , in this study. According to the previous study [22], the geometry of a crack could be known from heat tinting of the smooth, round, bar specimens ruptured at 1089 K. The crack depth, a , of the surface-notched specimens could also be known from the surface-crack length, l , and the crack geometry. The size of the surface-crack length and crack depth includes that of the notched part. The crack-growth rate was obtained by plotting crack depth, a , against time, t , and reading the slope of the plotted curve.

2.2. Estimation of the fractal dimensions of the grain-boundary fracture and of the grain boundaries

The smooth bar specimens, ruptured under one creep stress (108 MPa), were longitudinally sectioned in the

TABLE I Chemical composition of the cobalt-base HS-21 alloy used in this study (wt %)

Alloy	C	Cr	Ni	Fe	Mn	Mo	Si	S	P	B	Co
HS-21	0.27	26.71	2.37	0.09	0.64	5.42	0.59	0.007	< 0.005	0.003	Bal.

TABLE II The heat treatments, the mean value, D_{gb} , and the standard deviation, s_{gb} , of the fractal dimension of the grain boundaries, and the matrix hardness in the specimens of the HS-21 alloy

Heat treatments ^a	D_{gb} , s_{gb} (range of the fractal dimension)	Matrix hardness (Vickers hardness number, load = 4.9 N)
1523 K – 3.6 ks → WQ + 1089 K – 10.8 ks → AC	1.056, 0.0112 (1.039–1.066)	418
1523 K – 3.6 ks → AC + 1089 K – 10.8 ks → AC	1.107, 0.0194 (1.085–1.136)	407
1523 K – 3.6 ks – FC → 1383 K → WQ + 1089 K – 10.8 ks → AC	1.162, 0.0361 (1.117–1.211)	419
1523 K – 3.6 ks – FC → 1323 K → WQ + 1089 K – 10.8 ks → AC	1.241, 0.0325 (1.208–1.288)	405

^a WQ = water quenched; AC = air cooled; FC = furnace cooled.

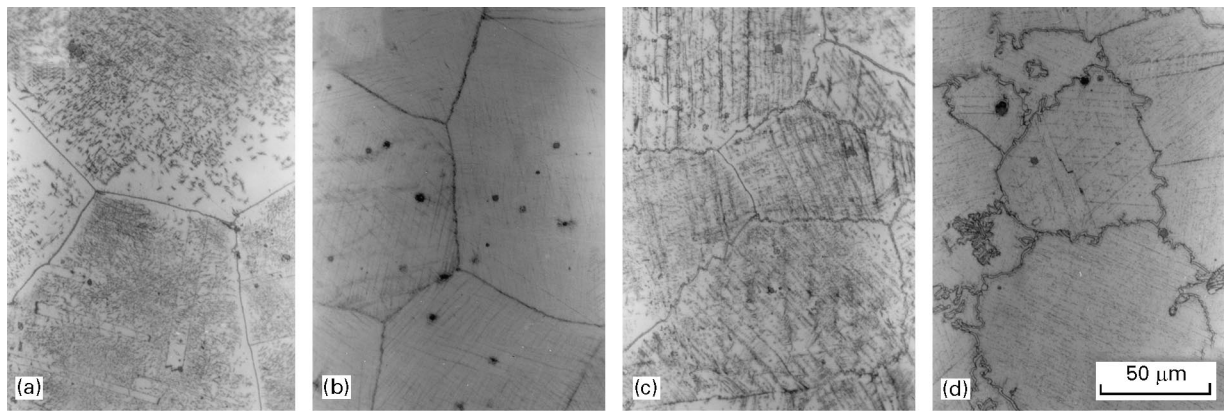


Figure 1 Optical micrographs of the heat-treated specimens of the HS-21 alloy. (a) $D_{gb} = 1.056$, (b) $D_{gb} = 1.107$, (c) $D_{gb} = 1.162$, (d) $D_{gb} = 1.241$.

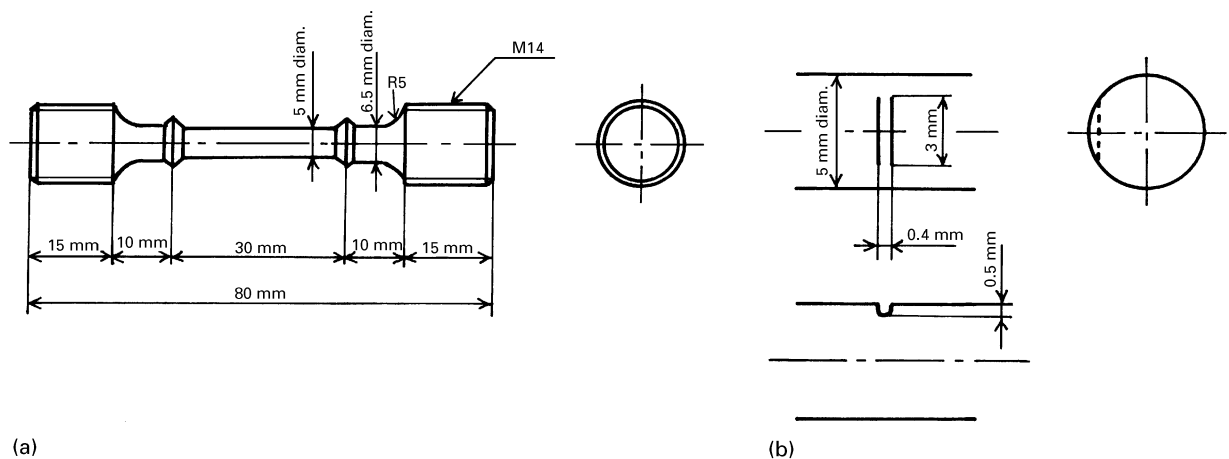


Figure 2 The geometry of the creep and creep-rupture specimens. (a) Smooth, round, bar specimen; (b) detail of the notched part of the surface-notched specimen.

plane involving the specimen axis, and the fracture surface profiles of five grain boundaries near the specimen axis were examined on each specimen. As well as in the measurement of the fractal dimension of the grain boundaries, D_{gb} , four representative photographs of the grain-boundary fracture surface profile were taken using an optical microscope at 400 times or using a scanning electron microscope at 2000 times. The mean value, D_f , and the standard deviation, s_f , of the fractal dimension of the grain-boundary fracture were estimated on five grain-boundary facets. The fractal dimension of the grain-boundary fracture in this study is the fractal dimension of the grain-boundary fracture surface profile in the two-dimensional section ($1 < D_f < 2$) estimated in the scale range, r , smaller than one grain-boundary length ($6.7 \times 10^{-7} \text{ m} \leq r \leq 1.3 \times 10^{-6} \text{ m}$, $5.2 \times 10^{-3} \leq r/d_g \leq 0.10$). The fractal dimension of the grain boundaries, D_{gb} , was also estimated in the similar scale range ($4.4 \times 10^{-3} \leq r/d_g \leq 8.8 \times 10^{-2}$).

Fig. 3 shows a schematic illustration of the box-counting method applied to the estimation of the fractal dimensions of the grain-boundary fracture and of the grain boundaries. The value of N_i is the number of boxes of the side length, r , which intersect the i th profile of the grain-boundary facet (or the i th profile of

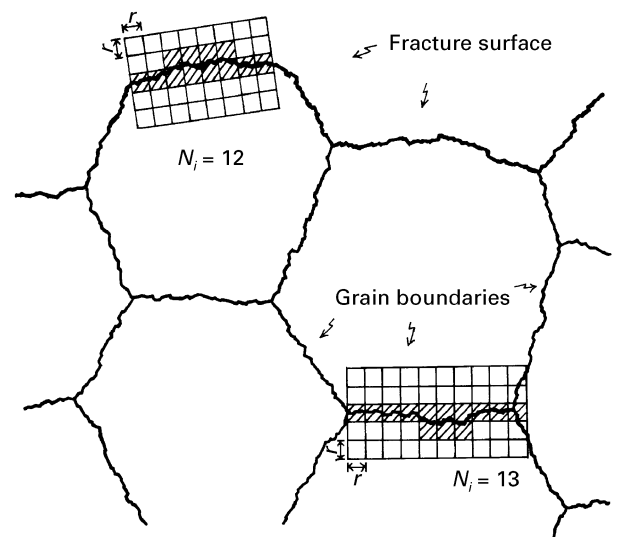


Figure 3 Schematic illustration of the box-counting method applied to the estimation of the fractal dimensions of the grain-boundary fracture and of the grain boundaries.

the grain-boundary surface). The length of the i th profile of the grain-boundary facet (or the i th profile of the grain-boundary surface). L_i , is given by a product

of N_i and r , such that $L_i = N_i r$. The length of the i th profile, L_i , is generally correlated to the scale length of the fractal analysis, r , by the following equation [14, 20, 21, 23]

$$L_i = L_{i0} r^{1-D_i} \quad (1)$$

where D_i is the fractal dimension of the i th profile and L_{i0} is a constant. The value of D_i can be obtained by fitting the experimental data of the i th profile to the above equation. The mean value, D , and the standard deviation, s , of the fractal dimension are given by

$$D = \sum_{i=1}^n D_i/n \quad (2)$$

$$s = \left[\sum_{i=1}^n (D_i - D)^2/(1-n) \right]^{1/2} \quad (3)$$

where $D = D_i$ and $s = s_f$ for the fractal dimension of the grain-boundary fracture, and $D = D_{gb}$ and $s = s_{gb}$ for the fractal dimension of the grain boundaries, and $n = 5$ in this study.

3. Results and discussion

3.1. Effects of grain-boundary configuration on the creep-rupture properties of smooth, round, bar specimens

Fig. 4 shows the rupture lives of the specimens in the HS-21 alloy at 1089 K. The rupture lives of the specimens tend to increase with increasing the mean value of the fractal dimension of the grain boundaries, D_{gb} . The increase in the value of D_{gb} leads to a large improvement of the rupture lives, especially under the higher stresses.

Fig. 5 shows the creep ductility of the specimens in the HS-21 alloy. Both the elongation and the reduction of area increase with increasing the value of D_{gb} . Thus, the improvement of the creep-rupture properties by serrated grain boundaries depend on the value of D_{gb} . The change in the creep-rupture properties with the value of D_{gb} can be explained by the initiation and growth of the grain-boundary crack during creep.

Fig. 6 shows examples of creep curves of the specimens of the HS-21 alloy tested under a stress of 108 MPa at 1089 K. The first grain-boundary crack seems to initiate on the specimen surface. The strain to crack initiation is also shown in this figure. As reported in the previous study [8], the creep curves of the HS-21 alloy have a long tertiary creep regime and a very short steady-state creep regime in the temperature range 1089–1311 K. The strain to the initiation of the surface crack is about 0.0083 in the specimen of $D_{gb} = 1.241$, whereas it is almost the same (about 0.0052–0.0056) in the specimens of D_{gb} less than 1.162.

Ashby and co-workers [9–11] discussed the influence of ledges, bumps, curvature, particles or precipitates on the grain-boundary sliding. Obviously, all these make the grain-boundary sliding difficult and increase the value of D_{gb} as well. According to the previous study [6, 7], serrated grain boundaries (with high D_{gb} value) are effective in retarding the initiation of the grain-boundary cracks controlled by grain-

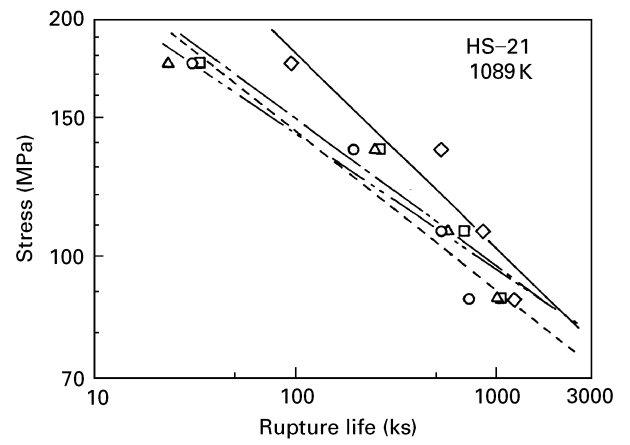


Figure 4 The rupture lives of the specimens in the HS-21 alloy at 1089 K, for the D_{gb} values: (—○—) 1.056, (—△—) 1.107, (—□—) 1.162, (—◇—) 1.241.

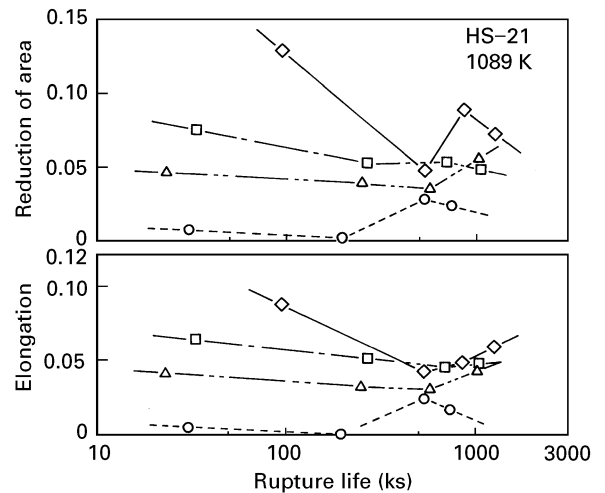


Figure 5 The creep ductility of the specimens in the HS-21 alloy at 1089 K. Values of D_{gb} given in Fig. 4.

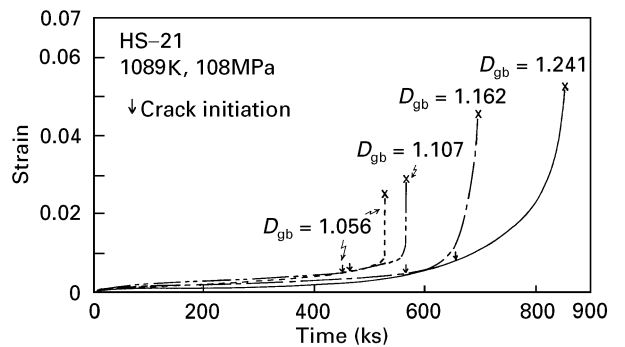


Figure 6 Creep curves of the specimens in the HS-21 alloy tested under a stress of 108 MPa at 1089 K.

boundary sliding, and the strain to crack initiation is larger in the specimen with serrated grain boundaries (with high D_{gb} value) compared with the one with straight grain boundaries (with low D_{gb} value). Therefore, the strain to crack initiation is the largest in the specimen of the largest D_{gb} value. However, there is almost no difference in the strain to crack initiation between the specimens of D_{gb} less than 1.162. This can be explained by boundary to boundary variation in

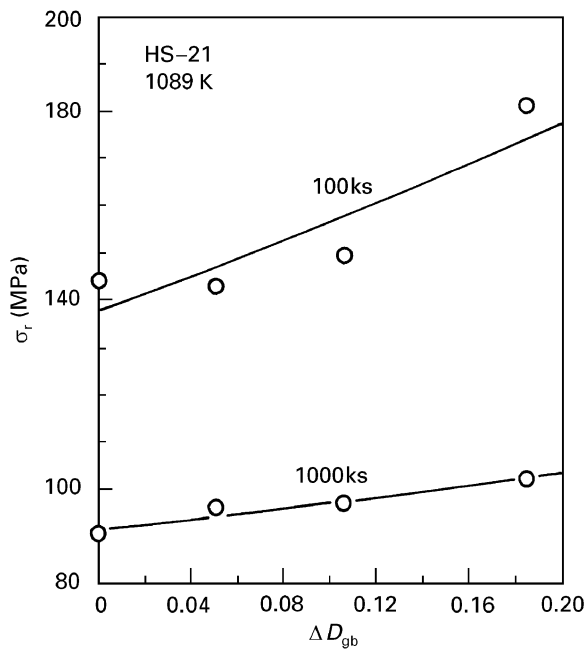


Figure 7 The relationship between the rupture strength, σ_r , and the increment in the mean value of the fractal dimension of the grain boundaries, ΔD_{gb} , in the HS-21 alloy at 1089 K ($\Delta D_{gb} = D_{gb} - 1.056$).

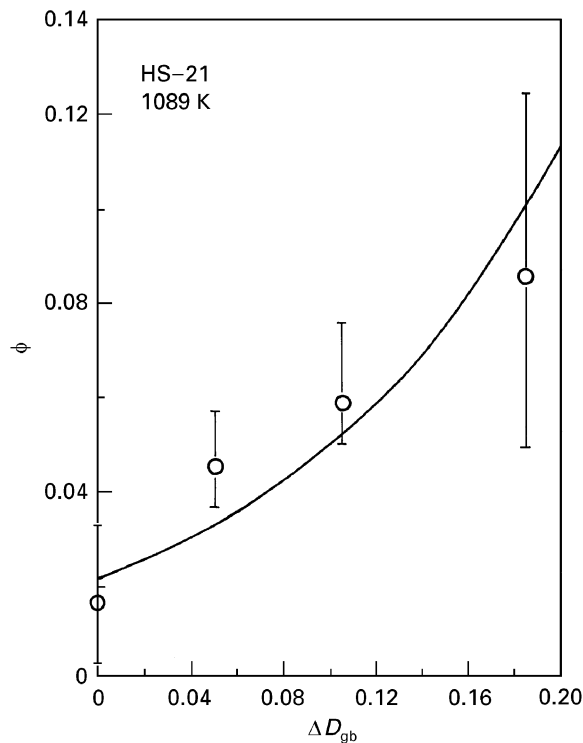


Figure 8 The relationship between the reduction of area, ϕ , and the increment in the mean value of the fractal dimension of the grain boundaries, ΔD_{gb} , in the HS-21 alloy at 1089 K ($\Delta D_{gb} = D_{gb} - 1.056$).

the fractal dimension of the grain boundaries in the specimens indicated by the standard deviation, s_{gb} , and the range of the fractal dimension (Table II). The crack initiation may be associated with the straighter grain boundaries (with the lower D_{gb} values) in the same specimen, although this was not confirmed in each specimen. As indicated by the value of s_{gb}

(Table II), the local variation of the fractal dimension seems to increase up to $D_{gb} = 1.162$. It is deduced that the lowest value of the fractal dimension in the specimen of $D_{gb} = 1.107$ and in the specimen of $D_{gb} = 1.162$ is close to that of the fractal dimension in the specimen of $D_{gb} = 1.056$, while there is a large difference in the lowest value of the fractal dimension between the specimen of $D_{gb} = 1.241$ and the specimen of $D_{gb} = 1.056$. Thus, the crack initiation may be associated not with the mean value of the fractal dimension of the grain boundaries, D_{gb} , but with the local variation of the fractal dimension in the specimens.

Fig. 7 shows the relationship between the rupture strength, σ_r , and the increment in the mean value of the fractal dimension of the grain boundaries, ΔD_{gb} , in the HS-21 alloy at 1089 K ($\Delta D_{gb} = D_{gb} - 1.056$). The value of σ_r increases with increase of the value of ΔD_{gb} , and the increase of σ_r with ΔD_{gb} is more remarkable at the shorter rupture life (100 ks). Fig. 8 shows the relationship between the reduction of area, ϕ , and the increment in the mean value of the fractal dimension of the grain boundaries, ΔD_{gb} , in the HS-21 alloy at 1089 K. The data points indicate the averaged values of the reduction of area over four creep stresses, and the ranges of the values are also shown in this figure. The reduction of area also increases with increasing value of D_{gb} . Similar results were obtained on the dependence of the elongation on the value of D_{gb} . These results can be correlated to the fracture mode in the specimens.

Fig. 9 shows the microstructures of the specimens in the HS-21 alloy ruptured under a stress of 108 MPa at 1089 K. The tensile direction is vertical in the optical micrographs. Long grain-boundary cracks are observed near the fracture surface in the specimens of the smaller values of D_{gb} (Fig. 9a, b), while shorter grain-boundary cracks can be seen in the specimens of the larger values of D_{gb} (Fig. 9c and d). The fracture surface of the specimen of the lowest D_{gb} value (1.056) is a typical brittle grain-boundary facet (Fig. 9e), whereas that of the specimen of the highest D_{gb} value (1.241) is a ductile grain-boundary fracture surface composed of ledges, steps and small dimples (Fig. 9h). The fracture surfaces of the specimens of the intermediate value of D_{gb} have a mixed microstructure of brittle and ductile fracture surfaces (Fig. 9f and g), which reflect the local variation of the fractal dimension of the grain boundaries in these specimens (Table II), and the area fraction of the ductile grain-boundary fracture surface is larger in the specimen of the higher D_{gb} value (Fig. 9g).

Table III shows the mean value, D_f , and the standard deviation, s_f , of the fractal dimension of the grain-boundary fracture and the range of the fractal dimension in the specimens of various D_{gb} values ruptured under a stress of 108 MPa at 1089 K. The value of D_f is close to that of D_{gb} , although it is a little higher than that of D_{gb} in the specimens of the lower D_{gb} values. The value of s_f is also close to that of s_{gb} in these specimens (Table II). The fracture appearance changes from a brittle grain-boundary fracture to a ductile one with increasing value of D_f (and the value of D_{gb}). It seems that the brittle to ductile transition in

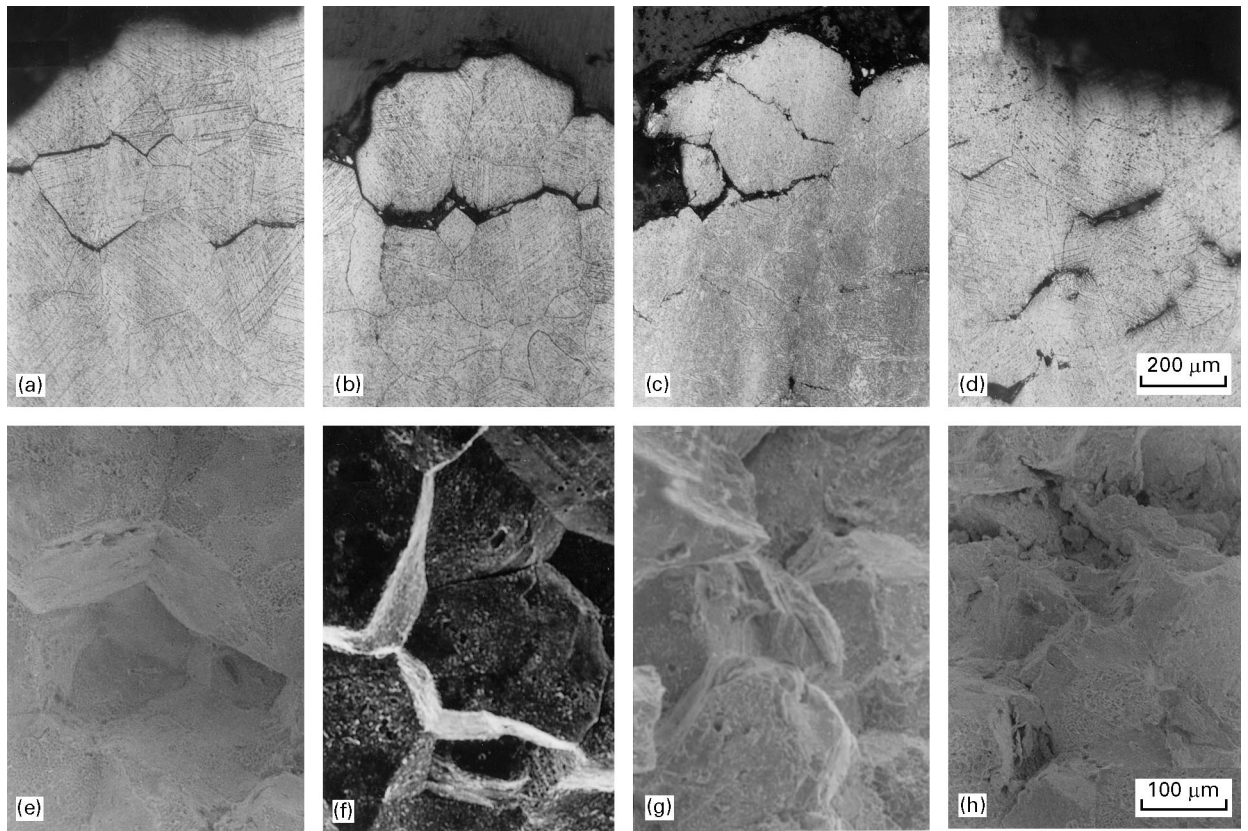


Figure 9 The microstructures of the specimens in the HS-21 alloy ruptured under a stress of 108 MPa at 1089 K. (a, e) $D_{gb} = 1.056$, (b, f) $D_{gb} = 1.107$, (c, g) $D_{gb} = 1.162$, (d, h) $D_{gb} = 1.241$.

TABLE III The mean value, D_f , and the standard deviation, s_f , of the fractal dimension of the grain-boundary fracture in the specimens of various D_{gb} values of the HS-21 alloy ruptured under the stress of 108 MPa at 1089 K

D_{gb}	D_f, s_f (range of the fractal dimension)
1.056	1.124, 0.0171 (1.099–1.139)
1.107	1.152, 0.0213 (1.126–1.179)
1.162	1.216, 0.0358 (1.182–1.271)
1.241	1.261, 0.0307 (1.212–1.289)

the fracture appearance occurs at the fractal dimension of the grain-boundary fracture of about 1.15–1.20 (at that of the grain boundaries of about 1.10–1.15). Further, the fracture appearance in the specimens was not changed with creep stress. Thus, both increases in the rupture strength and in the creep ductility with increasing D_{gb} value can be explained by an increase in the area fraction of the ductile grain-boundary fracture indicated by the value of D_f in the HS-21 alloy.

3.2. Creep crack growth in the surface-notched specimens

The growth of the grain-boundary crack during creep at 1089 K was investigated using three kinds of speci-

men with different D_{gb} values (1.056, 1.162 and 1.241). Fig. 10 shows the geometry of the grain-boundary cracks observed on the smooth, round, bar specimens of the HS-21 alloy ruptured at 1089 K. The product of specimen radius, $d/2$, and open angle, θ , gives the surface crack length ($l = d\theta/2$). The ratio of crack depth, a , to specimen diameter, d , i.e. a/d , increases with an increase in the angle θ (or the surface crack length, l), while the ratio R/d shows a constant value of 0.630 for all the specimens of the HS-21 alloy, where R is the radius of curvature at the crack front. The similar value of R/d (0.679) was obtained in the austenitic 21Cr–4Ni–9Mn steel [22]. Thus, if the value of l is known, the value of a/d can be estimated from the angle θ by using Fig. 10.

Fig. 11 shows the increment of the surface crack length, Δl , in the surface-notched specimens of the HS-21 alloy during creep at 1089 K. The period of the crack growth is longer in the specimen of the higher D_{gb} value under both gross stresses ($\sigma_g = 73.5, 88.2$ MPa), although the period of the crack growth in the surface-notched specimen is much less than that in the smooth bar specimen (Fig. 6). The time to crack initiation is not always shortest in the surface-notched specimen of the lowest D_{gb} value. This behaviour can be explained as follows. The grain boundary at which a crack is expected to initiate is not always located at the notch root where the stress concentration is the highest [22]. Further, as described in the previous section, there is a boundary-to-boundary variation in the fractal dimension of the grain boundaries. These factors may affect the initiation of the grain-boundary

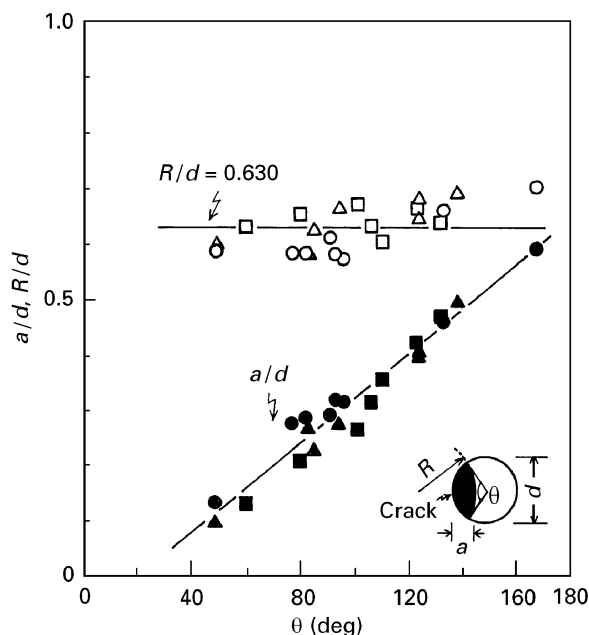


Figure 10 The geometry of the grain-boundary cracks observed on the smooth, round, bar specimens of the HS-21 alloy ruptured at 1089 K. D_{gb} : (Δ , \blacktriangle) 1.056, (\square , \blacksquare) 1.162, (\circ , \bullet) 1.241.

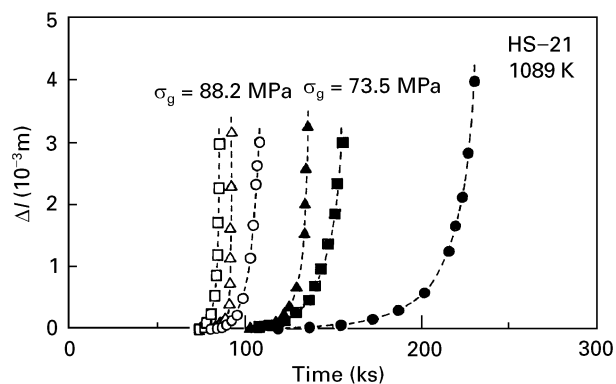


Figure 11 The increment of the surface crack length, Δl , with time in the surface-notched specimens of the HS-21 alloy during creep at 1089 K. D_{gb} : (Δ , \blacktriangle) 1.056, (\square , \blacksquare) 1.162, (\circ , \bullet) 1.241.

cracks. The change in the notch-opening displacement with time was also measured in the surface-notched specimens during creep at 1089 K, but it is not shown here, because the change in the notch-opening displacement was too small (less than 10^{-5} m) to identify through crack-growth experiments because of low creep ductility of the HS-21 alloy.

The crack depth, a , and the crack growth rate, da/dt , can be obtained from the results of the surface crack-growth experiments using the relationship between crack depth, a , and the angle θ (Fig. 10). Fig. 12 shows the relationship between the crack-growth rate, da/dt , and the maximum stress intensity factor, K_I , in the HS-21 alloy during creep at 1089 K. The K_I values for the surface-notched specimens is given using variables shown in Fig. 12 by Kiuchi *et al.*'s equation [24]

$$K_I = \sigma_g (\pi a)^{1/2} \{ 1.12 + (0.30/d - 1.85/R)a - [6.63/d^2 - 10.25/(Rd)]a^2 + [23.13/d^3 - 18.75/(Rd^2)]a^3 \} \quad (4)$$

$(a/d \leq 0.4 \text{ and } R \geq 0.5d)$

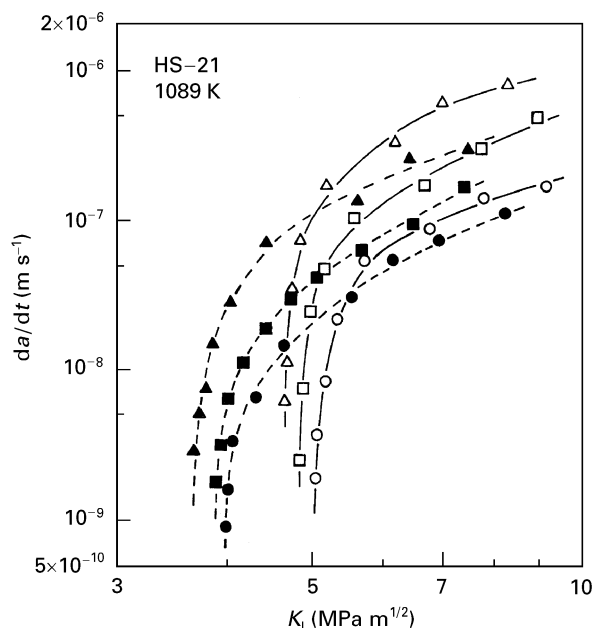


Figure 12 The relationship between the crack-growth rate, da/dt , and the maximum stress intensity factor, K_I , in the HS-21 alloy during creep at 1089 K, for $\sigma_g =$ (\blacktriangle , \blacksquare , \bullet) 73.5 MPa, and (Δ , \square , \circ) 88.2 MPa. D_{gb} : (\blacktriangle , Δ) 1.056, (\blacksquare , \square) 1.162, (\bullet , \circ) 1.241.

where σ_g is the gross stress. The crack-growth rate (da/dt) is lower in the specimens of the higher D_{gb} values, but the difference in the value of da/dt was larger at the smaller K_I values (i.e. in the early stage of the crack growth) between these specimens. The threshold stress intensity factor for crack growth, below which the crack does not grow, is higher in the specimens of the higher D_{gb} values, although it seems to be a little higher under the higher gross stress (88.2 MPa). As described in the previous section, the mean value of the fractal dimension of the grain-boundary fracture, D_f , is close to that of the fractal dimension of the grain boundaries, D_{gb} , in the smooth, round, bar specimens. Therefore, the crack growth in the microscopic scale may be controlled by the grain-boundary microstructures also in the surface-notched specimens.

Kitagawa and Yuuki [25] have reported that the K_I value of a deflected crack depends on the ratio of the deflected portion of the total length of the crack, and that K_I decreases with increasing proportion of the deflected part of the crack, even for the same projected crack length on the original crack trajectory. Ishida [26] has shown that the K_I value of a serrated crack normal to the tensile axis is almost constant, independent of the number of serrations. Suresh [27] has also obtained similar results and has shown that the circumvention of the crack path by serrated grain boundaries may retard the crack growth by lengthening the crack path. The grain-boundary crack is deflected not only at ledges, bumps, steps or precipitates on grain boundaries but also at grain-boundary triple junctions. The effect of grain-boundary configuration may be reduced by the crack deflection at grain-boundary triple junctions when the length of the grain-boundary crack is more than several grain diameters. Similar results were obtained on the austenitic 21Cr-4Ni-9Mn heat-resisting steel [22]. As shown by

the value of s_{gb} , there may be a local variation of the grain-boundary microstructures at the crack front, but this variation may be averaged along the crack front during the crack-growth process. Therefore, the crack growth can be correlated to the mean value of the fractal dimension of the grain boundaries, D_{gb} , in each specimen.

4. Conclusions

The creep–rupture properties and the initiation and growth of the grain-boundary cracks were investigated using four kinds of specimen of the cobalt-base HS-21 alloy with various grain-boundary microstructures at 1089 K in air. The results obtained were summarized.

1. Both the rupture strength and the creep ductility increased with increasing mean value of the fractal dimension of the grain boundaries, D_{gb} , in the specimens. The strain to the initiation of the grain-boundary crack was largest in the specimen of the highest D_{gb} value (1.241), while the strain was much the same in the specimens of the D_{gb} value less than 1.162. This behaviour was explained by the local variation in the grain-boundary microstructures indicated by the standard deviation of the fractal dimension of the grain boundaries, s_{gb} .

2. The mean value D_f , and the standard deviation, s_f , of the fractal dimension of the grain-boundary fracture were, respectively, close to the values of D_{gb} and s_{gb} although the value of D_f was a little higher than the value of D_{gb} in the specimens of the lower D_{gb} values. The creep–rupture properties were correlated to the microstructural features of the grain boundaries and of the grain-boundary fracture surfaces indicated by these fractal dimensions, D_{gb} and D_f . The fracture appearance changed from a brittle grain-boundary fracture to a ductile one with increasing values of D_f and D_{gb} .

3. The growth rate of the grain-boundary crack was lower in the specimens of the higher D_{gb} values, because the creep crack growth was possibly controlled by the averaged grain-boundary microstructures at the crack front. The effect of the grain-boundary configuration on the crack growth was larger at the smaller K_I value (at the early stage of the crack growth). The threshold stress intensity factor for crack growth was higher in specimens with higher D_{gb} values.

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References

1. W. BETTERIDGE and A. W. FRANKLIN, *J. Inst. Metals* **85** (1956–57) 473.
2. M. YAMAZAKI, *J. Jpn Inst. Metals* **30** (1966) 1032.
3. V. LUPINC, “High Temperature Alloys for Gas Turbines 1982” (Reidel, Dordrecht, 1982) p. 395.
4. M. KOBAYASHI, O. MIYAGAWA, T. SAGA and D. FUJISHIRO, *J. Iron Steel Inst. Jpn* **58** (1972) 79.
5. M. YAMAMOTO, O. MIYAGAWA, M. KOBAYASHI and D. FUJISHIRO, *ibid.* **63** (1977) 1848.
6. M. TANAKA, O. MIYAGAWA, T. SAKAKI and D. FUJISHIRO, *ibid.* **65** (1979) 939.
7. M. TANAKA, O. MIYAGAWA, T. SAKAKI, H. IIZUKA, F. ASHIHARA and D. FUJISHIRO, *J. Mater. Sci.* **23** (1988) 621.
8. M. TANAKA, H. IIZUKA and F. ASHIHARA, *ibid.* **24** (1989) 1623.
9. M. F. ASHBY, R. RAJ and R. C. GIFKINS, *Scripta Metall.* **4** (1970) 737.
10. R. RAJ and M. F. ASHBY, *Metall. Trans.* **2** (1971) 1113.
11. F. W. CROSSMAN and M. F. ASHBY, *Acta Metall.* **23** (1975) 425.
12. B. B. MANDELROT, D. E. PASSOJA and A. J. PAULLAY, *Nature* **308** (1984) 721.
13. V. Y. MILMAN, N. A. STELMASHENKO and R. BLUMENFELD, *Progr. Mater. Sci.* **38** (1994) 425.
14. B. B. MANDELROT, “The Fractal Geometry of Nature,” translated by H. Hironaka (Nikkei Science, Tokyo, 1985) p. 25.
15. M. TANAKA, *J. Mater. Sci.* **27** (1992) 4717.
16. *Idem, ibid.* **28** (1993) 5753.
17. M. TANAKA and H. IIZUKA, *Z. Metallkde.* **82** (1991) 442.
18. A. M. GOKHALE, W. J. DRURY and S. MISHRA, ASTM STP 1203, edited by J. E. Masters and L. N. Gilbertson (American Society for Testing and Materials, Philadelphia, PA, 1993) p. 3.
19. W. J. DRURY and A. M. GOKHALE, *ibid.*, p. 58.
20. H. TAKAYASU, “Fractals in the Physical Sciences,” (Manchester University Press, Manchester, New York, 1990) p. 11.
21. S. ISHIMURA and S. ISHIMURA, “Fractal Mathematics,” (Tokyo Tosho, Tokyo, 1990) p. 246.
22. M. TANAKA, H. IIZUKA and F. ASHIHARA, *J. Mater. Sci.* **23** (1988) 3827.
23. R. H. DAUSKARDT, F. HAUBENSAK and R. O. RITCHIE, *Acta Metall.* **38** (1990) 142.
24. A. KIUCHI, M. AOKI, M. KOBAYASHI and K. IKEDA, *J. Iron Steel Inst. Jpn* **68** (1982) 1830.
25. H. KITAGAWA and R. YUUKI, *Trans. Jpn. Soc. Mech. Eng.* **41-346** (1975) 1641.
26. M. ISHIDA, *Trans. Jpn. Soc. Mech. Eng.* **45-392** (1979) 306.
27. S. SURESH, *Metall. Trans.* **14A** (1983) 2375.

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